

5.6.2 Transformation Properties of Tensors

The last property of tensors that we need to consider is the manner in which they transform between different reference frames. This can be derived in a direct fashion that makes use of the fact that the tensors were defined in a way that is independent of the choice of reference frame, and hence tensors are geometrical objects (in the same way as four-vectors are) that have an existence in spacetime independent of any choice of reference frame. With that in mind, we can immediately write for the covariant components of a tensor T

$$T(\vec{d}, \vec{b}, \vec{c}, \dots) = T_{\mu\nu\alpha\dots} a^\mu b^\nu c^\alpha \dots = T_{\mu'\nu'\alpha'\dots} a^{\mu'} b^{\nu'} c^{\alpha'} \dots \quad (5.99)$$

By using $a^{\mu'} = \Lambda_{\nu}^{\mu'} a^{\nu}$ and similarly for the other vector components, this becomes

$$T_{\mu\nu\alpha\dots} a^{\mu} b^{\nu} c^{\alpha} \dots = T_{\mu'\nu'\alpha'\dots} \Lambda_{\mu}^{\mu'} \Lambda_{\nu}^{\nu'} \Lambda_{\alpha}^{\alpha'} \dots a^{\mu} b^{\nu} c^{\alpha} \dots \quad (5.100)$$

As the vectors $\vec{a}, \vec{b}, \vec{c} \dots$ are arbitrary, we have

$$T_{\mu\nu\alpha\dots} = \Lambda_{\mu}^{\mu'} \Lambda_{\nu}^{\nu'} \Lambda_{\alpha}^{\alpha'} \dots T_{\mu'\nu'\alpha'\dots} \quad (5.101)$$

In other words, the transformation is carried out in the same fashion as we have seen for the single index case (i.e. for the components of vectors). In a similar way (by use of $g^{\mu\nu}$ to raise indices), we can show for the contravariant components that

$$T^{\mu\nu\alpha\dots} = \Lambda_{\mu'}^{\mu} \Lambda_{\nu'}^{\nu} \Lambda_{\alpha'}^{\alpha} \dots T^{\mu'\nu'\alpha'\dots} \quad (5.102)$$

The results Eq. (5.101) and Eq. (5.102), and a corresponding result for mixed components of the tensor \mathbf{T} can be used as a test to see whether or not a multi-indexed quantity is, in fact, a tensor. We shall see how this can be implemented in the case of the Faraday tensor used to describe the electromagnetic field.